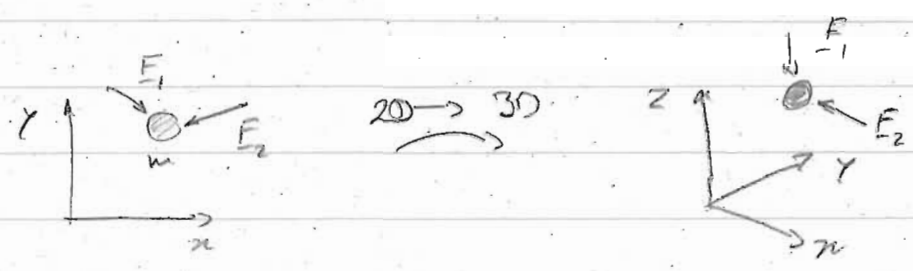


From 2D to 3D

We want to know the Newton-Euler Equations in 3D for a Rigid Body!

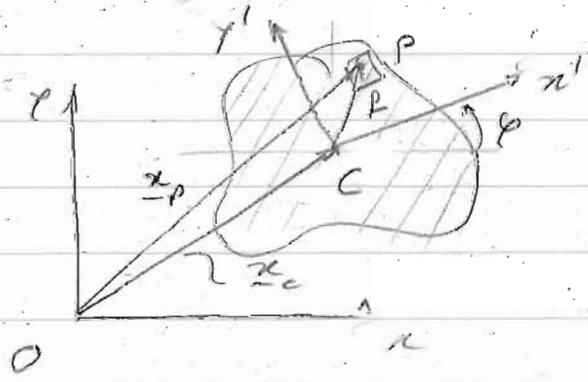
The Newton part is easy, just a point mass



$$\begin{aligned} \sum F_{ix} &= m\ddot{x} \\ \sum F_{iy} &= m\ddot{y} \end{aligned} \Rightarrow \sum \underline{F}_i = m\underline{\ddot{x}}$$

$$\begin{aligned} \sum F_{ip} &= m\ddot{x} \\ \sum F_{iq} &= m\ddot{y} \\ \sum F_{ir} &= m\ddot{z} \end{aligned} \Rightarrow \sum \underline{F}_i = m\underline{\ddot{x}}$$

The Euler part is about Rotation. Recap 2D:



Position cm  $(x_c, y_c)$   
 Orientation  $\varphi$

C-x'y' coordinate system is fixed to the rigid body → algebraic components of P are constant in this coordinate system.

Finding P in the frame of reference, the O-xy coordinate system

$$\underline{x}_p = \underline{x}_c + \underline{x}_{p/c}$$

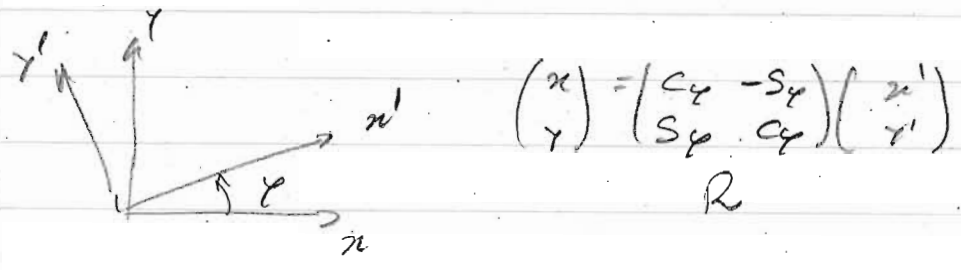
(2)

But  $\underline{x}'_{pc}$  is constant, rigid body

$$\underline{x}_p = \underline{x}_c + R(\varphi) \underline{x}'_{pc}$$

Any point on the rigid body can be found from a translation  $\underline{x}_c$  and a rotation  $\varphi$ .

The Rotation matrix in 2D is  $R(\varphi) = \begin{pmatrix} c_\varphi & -s_\varphi \\ s_\varphi & c_\varphi \end{pmatrix}$



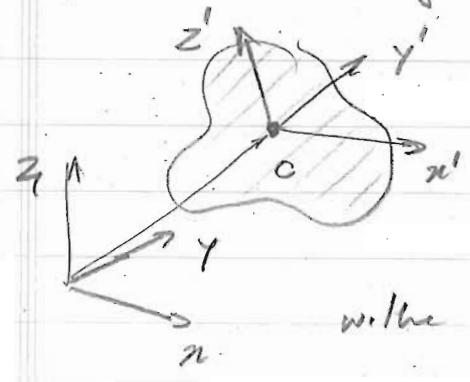
After adding up all point masses in the Newton eqn. for a 2D point mass

$$\left. \begin{aligned} \sum F_x &= m \ddot{x}_c \\ \sum F_y &= m \ddot{y}_c \\ \sum M_c &= J_c \ddot{\varphi} \end{aligned} \right\} \begin{array}{l} \text{Newton} \\ \text{Euler} \end{array}$$

(2\*)

with  $J_c = \int_V (x'^2 + y'^2) dm$  Mass moment of Inertia

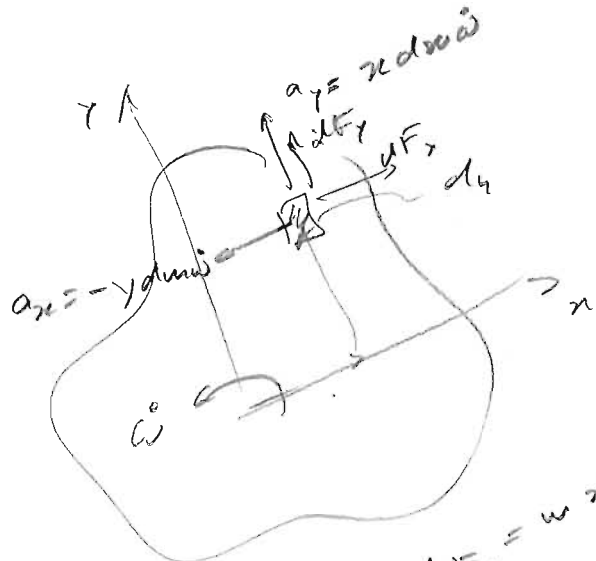
In 3D we get:



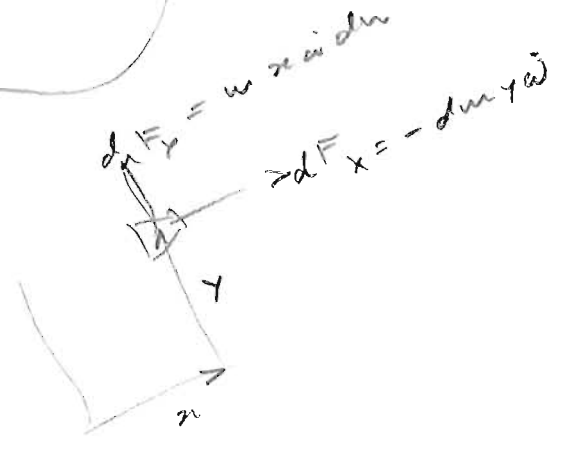
$$\underline{\Sigma F} = m \underline{\ddot{x}}_c \quad \text{Newton}$$

$$\underline{\Sigma M}_c = J_c \underline{\ddot{\omega}} + \underline{\omega} \times (J_c \underline{\omega}) \quad \text{Euler}$$

with the angular velocity vector  $\underline{\omega} = \begin{pmatrix} \omega_x \\ \omega_y \\ \omega_z \end{pmatrix}$



$f = ma$

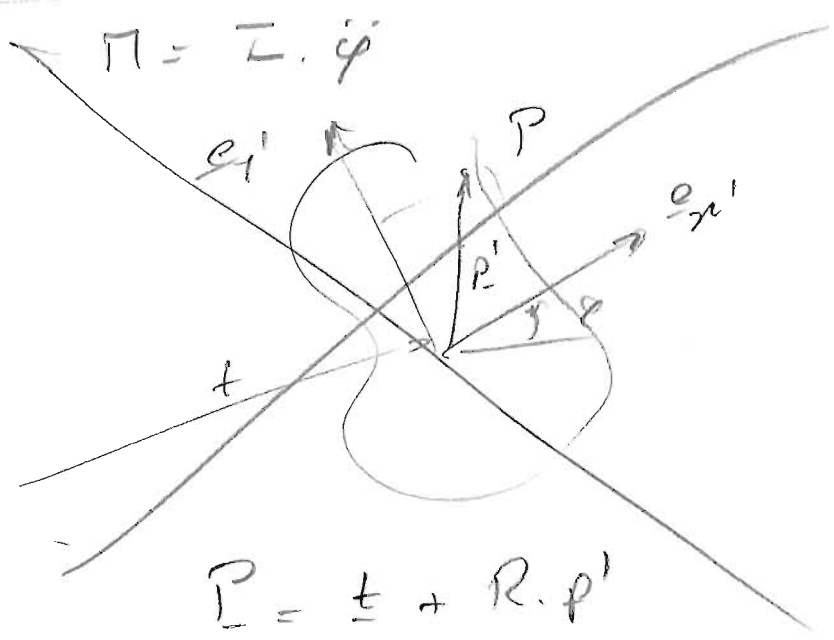


$$dM = x dF_y - y dF_x$$

$$= x x \dot{\omega} dm + y (-y \dot{\omega} dm)$$

$$dM = (x^2 + y^2) dm \dot{\omega}$$

$$M = \int (x^2 + y^2) dm \dot{\omega}$$



and the mass moment of inertia matrix

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$$J_c = \begin{pmatrix} J_{xx} & J_{xy} & J_{xz} \\ J_{yx} & J_{yy} & J_{yz} \\ J_{zx} & J_{zy} & J_{zz} \end{pmatrix} \quad \text{w/1} \quad \begin{aligned} J_{xx} &= \int (y^2 + z^2) dm & \text{22} \\ J_{xy} &= - \int xy dm & \text{222} \end{aligned}$$

Or along the principal axis, you get rid of the off-diagonal terms:

$$J_c = \begin{pmatrix} J_x & 0 & 0 \\ 0 & J_y & 0 \\ 0 & 0 & J_z \end{pmatrix}$$

( $J_x, J_y$  and  $J_z$  are the eigenvalues of  $J_c$ )

Back to Euler  $\Sigma \underline{M}_c = J_c \underline{\ddot{\omega}} + \underline{\omega} \times (J_c \underline{\omega})$

If the body rotates then  $J_c$  is not constant!  
Better to express Euler in body fixed coordinate system C-axes'

$$\Sigma \underline{M}'_c = J'_c \underline{\ddot{\omega}}' + \underline{\omega}' \times (J'_c \underline{\omega}')$$

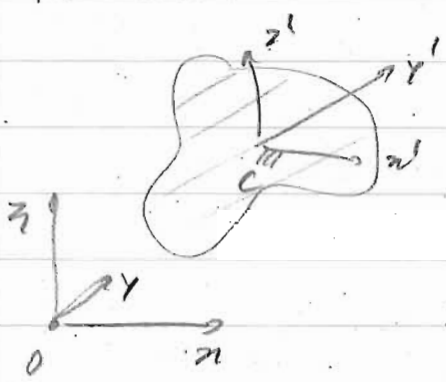
$\underline{\omega}'$  the angular velocity in the C-axes' coord system (local)

$\underline{M}'_c$  the moment in the local coord system

$J'_c$  is now a constant matrix.

Summarize Newton-Euler in 3D

Newton-Euler eqns of motion for a rigid body in 3D:



$$\begin{aligned} \Sigma F_x &= m \ddot{x}_c \\ \Sigma F_y &= m \ddot{y}_c \\ \Sigma F_z &= m \ddot{z}_c \end{aligned}$$

$$\begin{aligned} \Sigma M'_x &= J'_x \dot{\omega}'_x - (J'_y - J'_z) \omega'_y \omega'_z \\ \Sigma M'_y &= J'_y \dot{\omega}'_y - (J'_z - J'_x) \omega'_z \omega'_x \\ \Sigma M'_z &= J'_z \dot{\omega}'_z - (J'_x - J'_y) \omega'_x \omega'_y \end{aligned}$$

C, x', y', z' body fixed coord system with principal axes

$$J' = \begin{pmatrix} J'_x & 0 & 0 \\ 0 & J'_y & 0 \\ 0 & 0 & J'_z \end{pmatrix}$$

These 3D Euler equations look strange and complex, but... 3D motion of rigid bodies can be strange and complex

For example: the stability of torque free 3D rotation of a rigid body (think of a satellite, or tennis racket or a vanguard control)

Assume we start a rigid body with one big angular velocity and the other two zero  $\omega = (\omega_x, 0, 0)$ . Will this stay this way or will the body start rotating along one of the other axes in other words is this rotation stable??

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Everything in a body fixed coord system with principle axis along 1, 2, 3 axis in centre of mass.

Assume an initial angular velocity  $\underline{\omega}_0 = (\omega_1, 0, 0)$

The inertia matrix is  $J_c = \text{diag}(J_1, J_2, J_3)$

Target force  $\Sigma \underline{M} = \underline{0}$

Euler in local coord system:  $\Sigma \underline{M} = J \underline{\dot{\omega}} + \underline{\omega} \times (J \underline{\omega})$

We assume small changes from the initial conf.

$$\Sigma \underline{M} = \underline{M}_0 + \Delta \underline{M}, \quad \underline{\omega} = \underline{\omega}_0 + \Delta \underline{\omega} \rightarrow \underline{\dot{\omega}} = \underline{\dot{\omega}}_0 + \Delta \underline{\dot{\omega}}$$

$$\underline{M}_0 + \Delta \underline{M} = J \cdot (\underline{\dot{\omega}}_0 + \Delta \underline{\dot{\omega}}) + (\underline{\omega}_0 + \Delta \underline{\omega}) \times (J \cdot (\underline{\omega}_0 + \Delta \underline{\omega}))$$

$$\underline{M}_0 + \Delta \underline{M} = J \underline{\dot{\omega}}_0 + J \Delta \underline{\dot{\omega}} + \underline{\omega}_0 \times J \underline{\omega}_0 + \underline{\omega}_0 \times J \Delta \underline{\omega} + \Delta \underline{\omega} \times J \underline{\omega}_0 + \Delta \underline{\omega} \times J \Delta \underline{\omega}$$

$O(\Delta^2)$

$$\underline{M}_0 = J \underline{\dot{\omega}}_0 + \underline{\omega}_0 \times J \underline{\omega}_0 \rightarrow \underline{\dot{\omega}}_0 = \underline{0} \quad \{ \underline{\omega}_0 \text{ constant} \}$$

$$\Delta \underline{M} = J \Delta \underline{\dot{\omega}} + \underline{\omega}_0 \times J \Delta \underline{\omega} + \Delta \underline{\omega} \times J \underline{\omega}_0$$

$\Delta \underline{M} = \underline{0}$  We apply the entire Torques

$$\underline{0} = \begin{pmatrix} J_1 & & \\ & J_2 & \\ & & J_3 \end{pmatrix} \begin{pmatrix} \Delta \dot{\omega}_1 \\ \Delta \dot{\omega}_2 \\ \Delta \dot{\omega}_3 \end{pmatrix} + \begin{pmatrix} \omega_1 \\ 0 \\ 0 \end{pmatrix} \times \begin{pmatrix} J_1 \Delta \omega_1 \\ J_2 \Delta \omega_2 \\ J_3 \Delta \omega_3 \end{pmatrix} + \begin{pmatrix} \Delta \omega_1 \\ \Delta \omega_2 \\ \Delta \omega_3 \end{pmatrix} \times \begin{pmatrix} J_1 \omega_1 \\ 0 \\ 0 \end{pmatrix}$$
$$+ \begin{pmatrix} 0 \\ -\omega_1 J_3 \Delta \omega_3 \\ \omega_1 J_2 \Delta \omega_2 \end{pmatrix} + \begin{pmatrix} 0 \\ \Delta \omega_3 J_1 \omega_1 \\ -\Delta \omega_2 J_1 \omega_1 \end{pmatrix}$$

used here

$\underline{\omega} = (\omega, 0, 0)$   
 $\Delta \underline{\omega} = (\Delta \omega, \beta, \gamma)$   
 $J = (A, B, C)$

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$$\underline{0} = \begin{pmatrix} J_1 & 0 & 0 \\ 0 & J_2 & 0 \\ 0 & 0 & J_3 \end{pmatrix} \begin{pmatrix} \Delta \dot{\omega}_1 \\ \Delta \dot{\omega}_2 \\ \Delta \dot{\omega}_3 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & (J_1 - J_3)\omega_1 \\ 0 & (J_2 - J_1)\omega_1 & 0 \end{pmatrix} \begin{pmatrix} \Delta \omega_1 \\ \Delta \omega_2 \\ \Delta \omega_3 \end{pmatrix}$$

Assume a solution of the form  $\Delta \underline{\omega} = \underline{a} e^{\lambda t}$

This leads to the eigenvalue problem

$$\begin{pmatrix} J_1 \lambda & 0 & 0 \\ 0 & J_2 \lambda & (J_1 - J_3)\omega_1 \\ 0 & (J_2 - J_1)\omega_1 & J_3 \lambda \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix} \cdot e^{\lambda t} = \underline{0} \rightarrow$$

$$\text{Det} = 0 \rightarrow J_1 J_2 J_3 \lambda^3 + 0 + 0 - 0 - (J_2 - J_1)\omega_1 (J_1 - J_3)\omega_1 J_1 \lambda = 0$$

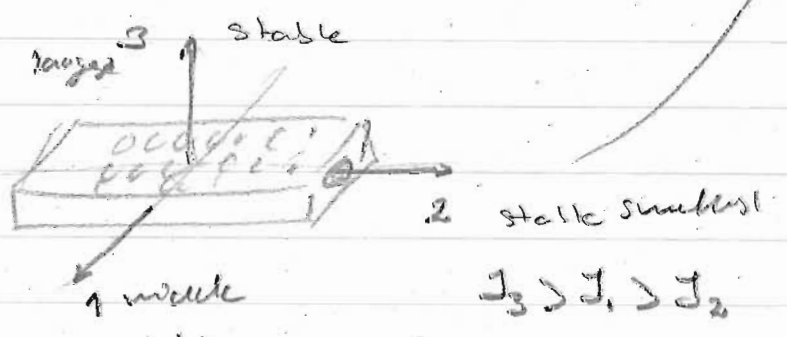
$$J_1 \lambda \{ J_2 J_3 \lambda^2 - (J_1 - J_3)(J_2 - J_1)\omega_1^2 \} = 0$$

$$\rightarrow \lambda_1 = 0 \quad \text{and} \quad \lambda_{2,3} = \sqrt{\frac{(J_1 - J_3)(J_2 - J_1)}{J_2 J_3}} \omega_1$$

Unstable  $\lambda_{2,3} > 0 \Rightarrow J_1 > J_3 \text{ and } J_2 > J_1 \Rightarrow J_2 > J_1 > J_3$   
 $\text{or } J_3 > J_1 \text{ and } J_1 > J_2 \Rightarrow J_3 > J_1 > J_2$

Cone-like rotation about the intermediate axis is unstable!

Poinsot's Case:



Tennis Racket

